

Modular theory in Algebraic Quantum Field Theory

Half-sided Modular Inclusions, Standard Pairs
and beyond

Ian Koot

RSME Congress of Young Researchers 2025 January 16, 2025

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- ω : expected value for observables.

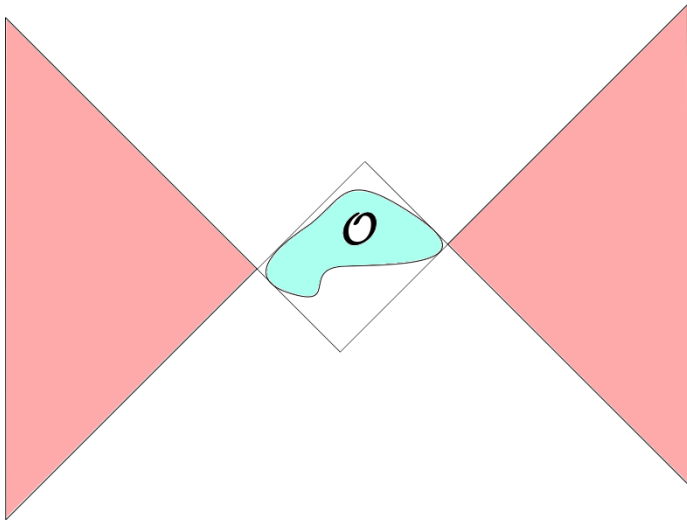
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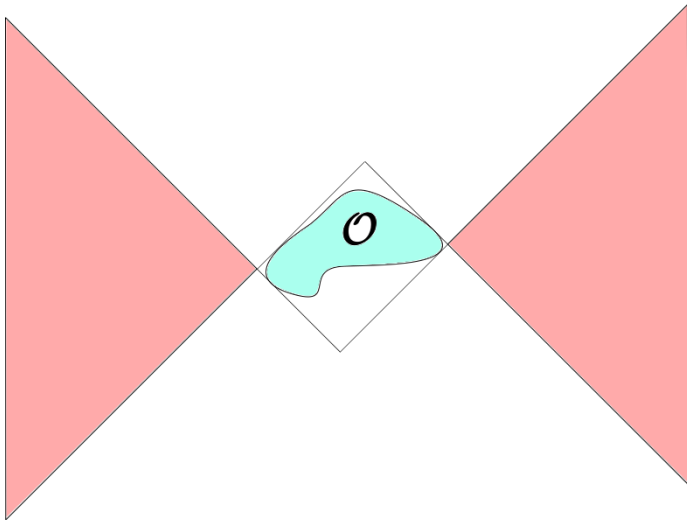
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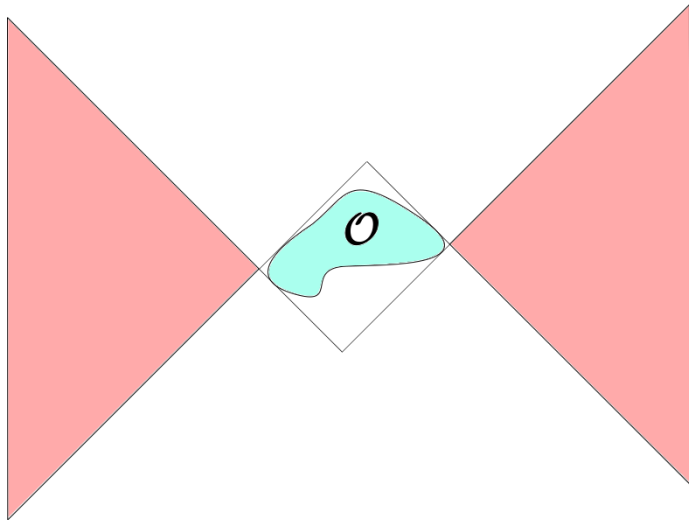
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Where are the 'fields'?

Classically: field is a **value** at each point in spacetime.

Quantum: field is an **observable** at each point in spacetime.

But these are usually singular/unbounded!

→ we 'smear' them

Let $\mathcal{A} \subset B(\mathcal{H})$ von Neumann Algebra, $\Omega \in \mathcal{H}$:

$$\text{cyclic: } \overline{\mathcal{A}\Omega} = \mathcal{H}, \quad \text{separating: } A\Omega = 0 \Rightarrow A = 0$$

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If Ω is **standard** for \mathcal{A} (i.e. cyclic and separating) one defines

$$S_{\mathcal{A},\Omega} : \mathcal{A}\Omega \rightarrow \mathcal{A}\Omega, \quad A\Omega \mapsto A^*\Omega$$

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Looks innocent, is usually unbounded, with $\overline{S_{\mathcal{A},\Omega}} = (S_{\mathcal{A},\Omega})^*$. Polar decomposition (with abuse of notation)

$$S_{\mathcal{A},\Omega} = J_{\mathcal{A},\Omega} \Delta_{\mathcal{A},\Omega}^{\frac{1}{2}}$$

These objects satisfy properties which *a priori* are not obvious at all:

$$\Delta^{it} \Omega = \Omega, \quad \Delta^{it} \mathcal{A} \Delta^{-it} = \mathcal{A}, \quad J \mathcal{A} J = \mathcal{A}'$$

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Example: matrix case

Let $\mathcal{H} := M_n(\mathbb{C})$ with $\langle A, B \rangle := \text{tr}(A^* B)$.

We take $\Omega = \rho^{\frac{1}{2}}$ (from $A \mapsto \text{tr}(\rho A)$)

Writing $L[A] : B \mapsto AB$ we have $\mathcal{A} := L[M_n(\mathbb{C})] \subset B(M_n(\mathbb{C}))$. Then $\mathcal{A}' = R[M_n(\mathbb{C})]$ (right multipliers).

The modular objects now are:

$$J(A) = A^*, \quad \Delta^{it}(A) = \rho^{it} A \rho^{-it}$$

so that indeed

$$J \Delta^{\frac{1}{2}}(L[A] \rho^{\frac{1}{2}}) = J(\rho^{\frac{1}{2}}(A \rho^{\frac{1}{2}}) \rho^{-\frac{1}{2}}) = A^* \rho^{\frac{1}{2}} = L[A]^* \rho^{\frac{1}{2}}$$

Theorem (Takesaki 1970)

Let $(\mathcal{H}, \mathcal{A}(\cdot), U, \Omega)$ be a QFT. Then the state $\langle \Omega, \cdot \Omega \rangle$ on $\mathcal{A}(\mathcal{O})$ is at “temperature” β w.r.t. the time evolution $U((t, 0))$ if and only if

$$\Delta_{\mathcal{A}(\mathcal{O}), \Omega}^{it} = U \left(\left(-\frac{\beta t}{2}, 0 \right) \right)$$

Theorem (Borchers 1992)

Let $(\mathcal{H}, \mathcal{A}(\cdot), U, \Omega)$ be a QFT, such that U is a **positive energy representation** and Ω is a **vacuum vector** for U . Then

$$\Delta_{\mathcal{A}(W), \Omega}^{it} U(\vec{x}) \Delta_{\mathcal{A}(W), \Omega}^{-it} = U(\Lambda_{2\pi t} \vec{x})$$

In Summary:

- QFT: von Neumann Algebras $\mathcal{A}(\mathcal{O})$ with **inclusion** and **commutation** relations + **unitary representation** U of \mathbb{R}^2 + a state ω (sometimes given by a vector Ω).

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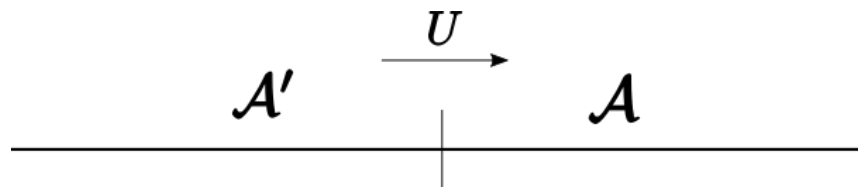
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- QFT: von Neumann Algebras $\mathcal{A}(\mathcal{O})$ with **inclusion** and **commutation** relations + **unitary representation** U of \mathbb{R}^2 + a state ω (sometimes given by a vector Ω).
- Modular theory: **von Neumann Algebra** and a **standard vector** \Rightarrow the **modular group** Δ^{it} and the **modular conjugation** J ‘for free’.
- We’re interested in the interaction between QFT and modular theory.

A simpler situation: suppose $\mathcal{A} \subset B(\mathcal{H})$ VNA, and $U : \mathbb{R} \rightarrow \mathcal{U}(\mathcal{H})$ *positively generated*. Geometrical assumption:

$$U(t)\mathcal{A}U(-t) \subset \mathcal{A} \quad \text{for } t \geq 0$$

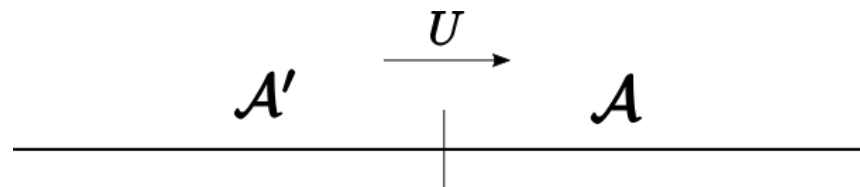
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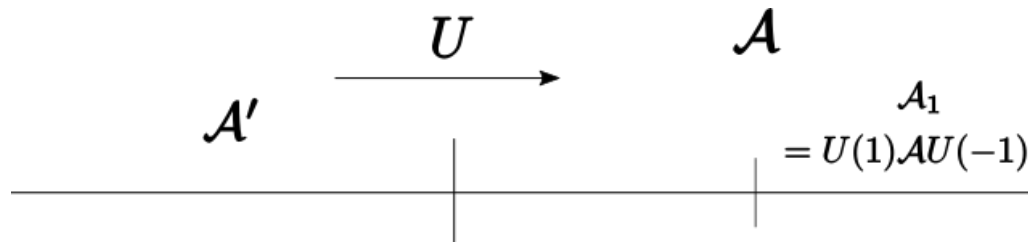
Theorem (Borchers 1992)

If (\mathcal{A}, U) as above, and Ω is standard for \mathcal{A} and invariant for U , then

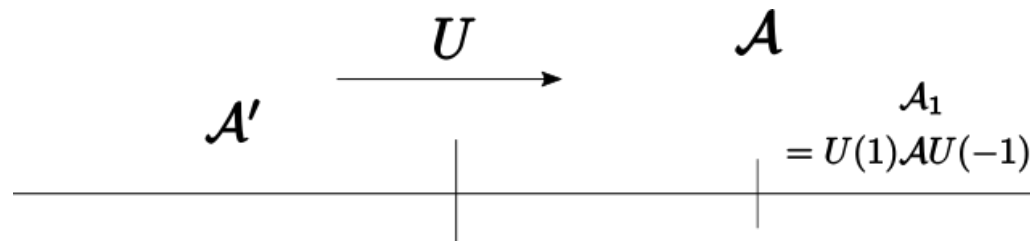
$$\Delta_{\mathcal{A}, \Omega}^{it} U(s) \Delta_{\mathcal{A}, \Omega}^{-it} = U(e^{-2\pi t} s)$$

Half-sided Modular Inclusions

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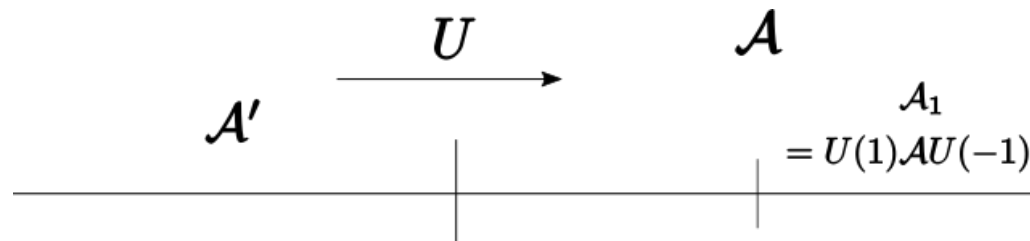


Then \mathcal{A}_1 satisfies a property only concerning the modular group:

$$\Delta_{\mathcal{A},\Omega}^{it} \mathcal{A}_1 \Delta_{\mathcal{A},\Omega}^{-it} = U(e^{-2\pi t}) \mathcal{A} U(-e^{-2\pi t}) \subset U(1) \mathcal{A} U(-1) = \mathcal{A}_1$$

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This is special! We call $\mathcal{A}_1 \subset \mathcal{A}$ a **half-sided modular inclusion** (w.r.t. Ω).

It is a totally non-obvious fact that *all* half-sided modular inclusions arise this way (Wiesbrock 1993, Araki-Zsido 2004):

$$\left\{ \mathcal{A}_1 \subset \mathcal{A}_2 \left| \begin{array}{l} \Omega \text{ std. for } \mathcal{A}_1 \\ \Delta_{\mathcal{A}, \Omega}^{it} \mathcal{A}_1 \Delta_{\mathcal{A}, \Omega}^{-it} \subset \mathcal{A}_1 \\ \text{for all } t \leq 0 \end{array} \right. \right\} \leftrightarrow \left\{ U : \mathbb{R} \rightarrow \mathcal{U}(\mathbb{R}) \left| \begin{array}{l} U \text{ pos. gen.} \\ U(t) \mathcal{A} U(-t) \subset \mathcal{A} \\ \text{for all } t \geq 0 \end{array} \right. \right\}$$

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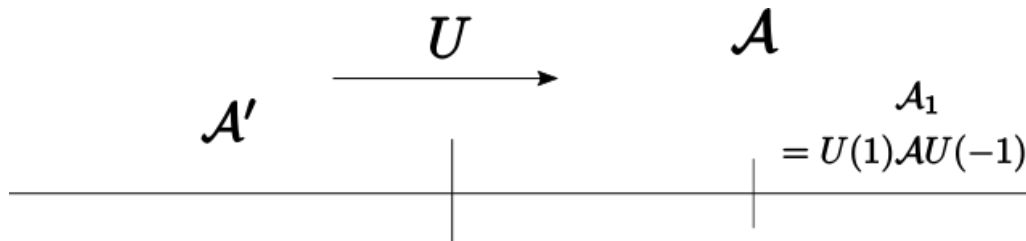
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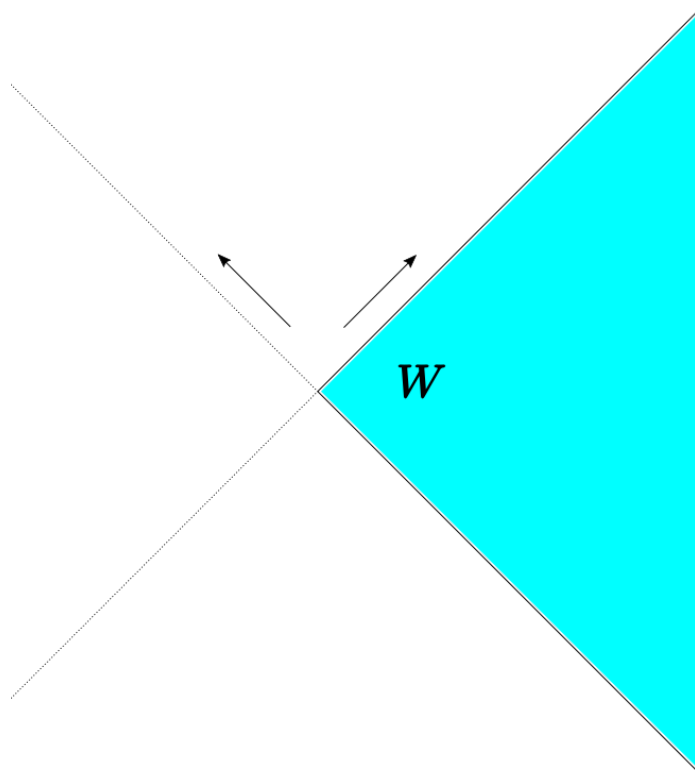
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From the perspective of the half-sided modular inclusions, there is a *hidden* geometry.

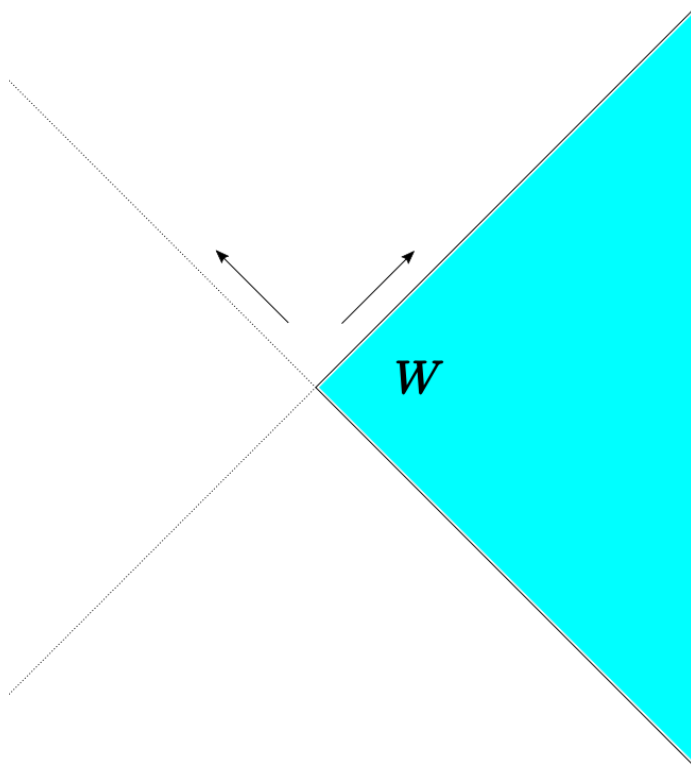




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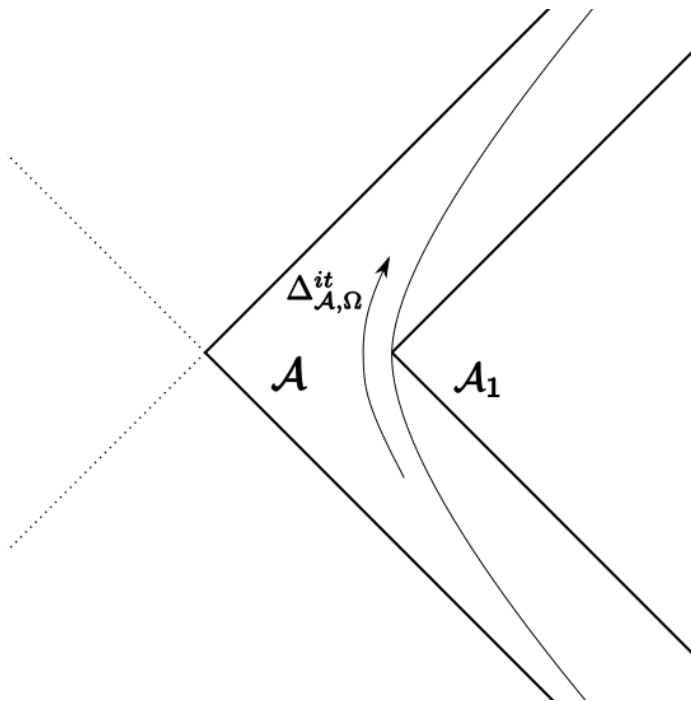
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In the above situation we have

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where Λ_s are the so called **Lorentz boosts**.



We are working on the following question:

Open Question

When can we reconstruct U from \mathcal{A} and \mathcal{A}_1 as in the 1-dimensional case?

The inclusion $\mathcal{A}_1 \subset \mathcal{A}$ is no longer a HSMI; but with an **extra** direction (corresponding to spatial translations) we can recognize two HSMI's along the lighttrays!

We're currently investigating when the two standard pair directions commute.

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- There is a correspondence between standard pairs (VNA + pos. gen. rep.) and half-sided modular inclusions (VNA + subVNA)
- We're researching a generalization to higher dimensions.

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